



Axisymmetric ($n = 0$) high-frequency MHD oscillations at the pedestal of JET plasmas

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Introduction

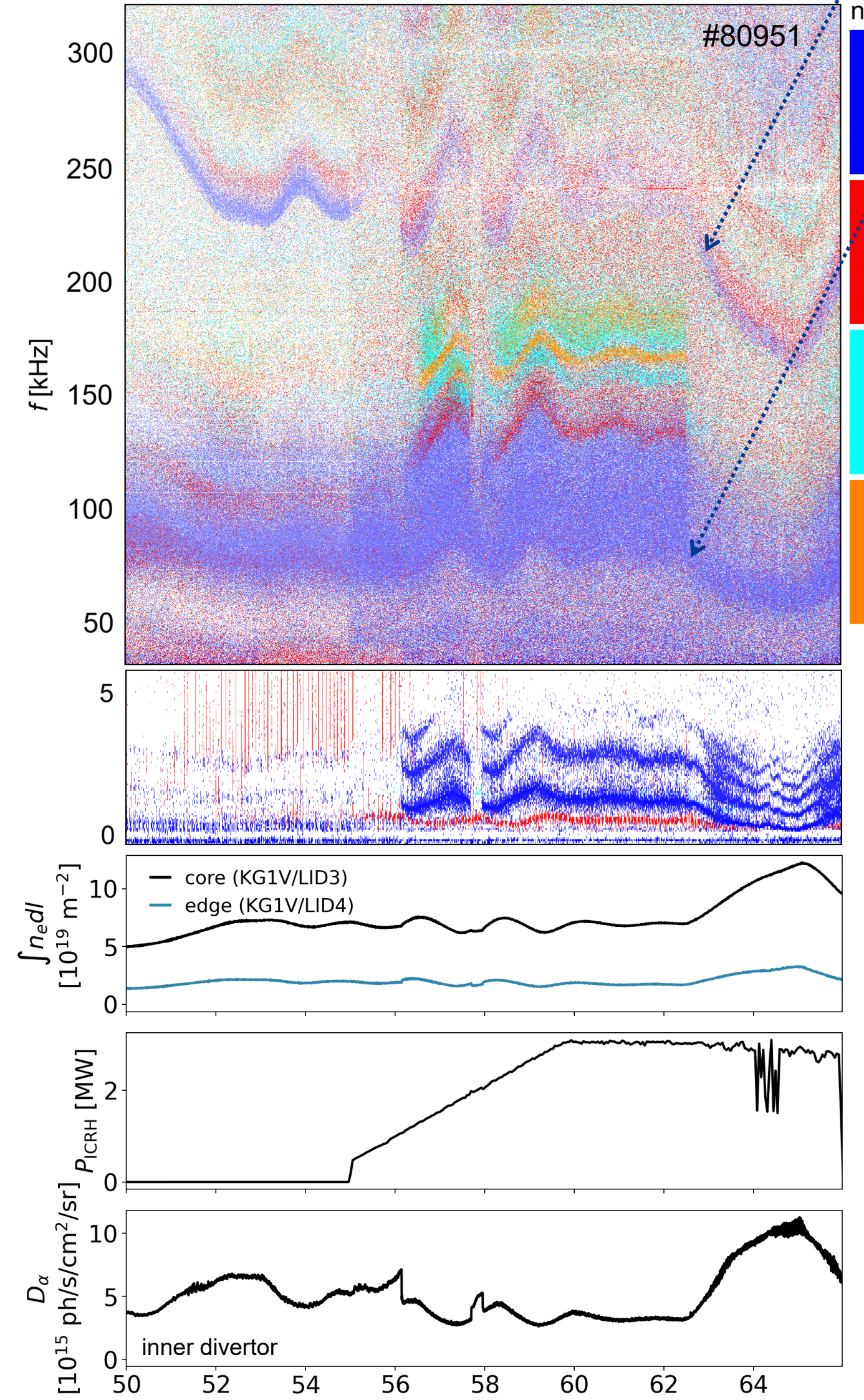
Long-lived axisymmetric ($n = 0$) quasi-coherent high frequency magnetic oscillations (HFOs, 50 – 500 kHz) in the TAE and EAE frequency range have been reported **close to the L-H transition** at JET [1] and AUG [2] tokamaks. Despite some models [3], their nature remains **elusive**.

HFOs are **not restricted to L-H transitions scenarios** and appear in a variety of plasma compositions (H, D, D-He³, D-T) with a wide range of plasma parameters and under different heating schemes (pure Ohmic, NBI and ICRH). Modes of Alfvénic nature with similar properties were observed on MAST [4] and COMPASS [5].

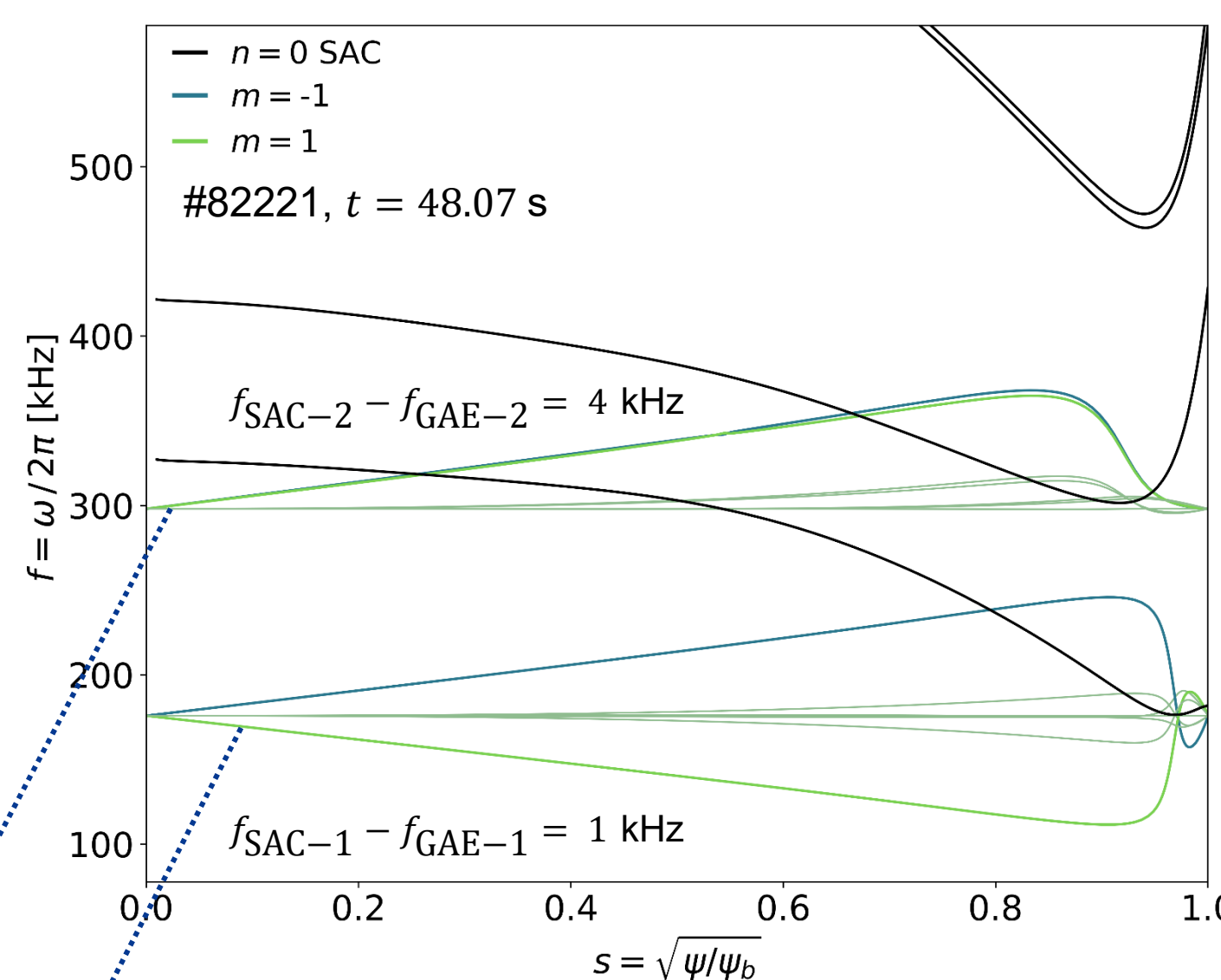
Here, we show that **axisymmetric HFOs** can be explained in light of **MHD theory** as **$n = 0$ Global Alfvén Eigenmodes (GAEs)**, whose frequencies we can robustly predict.

Experimental characterisation of HFOs

- HFOs appear **during L- and H-modes**;
- HFOs **persist for long periods** with relatively **low amplitude**;
- HFOs are **modulated by edge $n = 0$ $m = 1$ LFOs** emerging at the L-H transition;
 - M-mode [6] (or I-phase at AUG [2])
 - $f_{\text{M-mode}} \propto \frac{m}{r} v_A^{\text{poloidal}}$ [6]
- HFOs are present **before external heating power** is turned on, **while** it is on, and **prevail after** it is turned off;
- HFOs' signal is **stronger during Ohmic heating phase**;
- Generally, **NBI does not affect HFOs**;
- ICRH appears to suppress** the 2nd HFO;
- 2nd HFO is considerably **attenuated during the H-mode**;
- Both $n = 0$ HFOs are often accompanied by analogous $n = 1$ (and faint $n = 2$) bands.



Axisymmetric shear Alfvén waves and continuum



Only a class of shear Alfvén Eigenmodes with $n = 0$ can exist: the **Global Alfvén Eigenmodes**. They arise due to and near a **local extremum of the shear-Alfvén continuum (SAC)** produced by **radial profiles** [7].

Cylindrical limit $n = 0$ SAC: $\omega^2 \approx \frac{B^2}{\mu_0 \rho} \frac{m^2}{q^2 R^2}$

Near the edge, $q \uparrow$ and $\rho \downarrow \Rightarrow \omega^2$ minimum

MISHKA [8] finds **two GAEs below the minimum** of the first two branches of the SAC computed by CSMISH [9] (left plot).

Analytic expression for the two lowest branches of the coupled SAC

Motivation: $f_{\text{SAC}} \approx f_{\text{GAE}}$

Incompressible SAC $\equiv$ set of eigenvalues ω^2 that yields non-trivial solutions for [10]

$$\left[\frac{\omega^2 \rho |\nabla \Psi|^2}{B^2} + \mathbf{B} \cdot \nabla \left(\frac{|\nabla \Psi|^2}{B^2} \mathbf{B} \cdot \nabla \right) \right] \xi^A = 0$$

Analytic local magnetic equilibrium model [11]:

$$\Psi(r, \theta) = \Psi_b S_0 r^2 [\Theta_0(\theta) + \varepsilon r \Theta_1(\theta)]$$

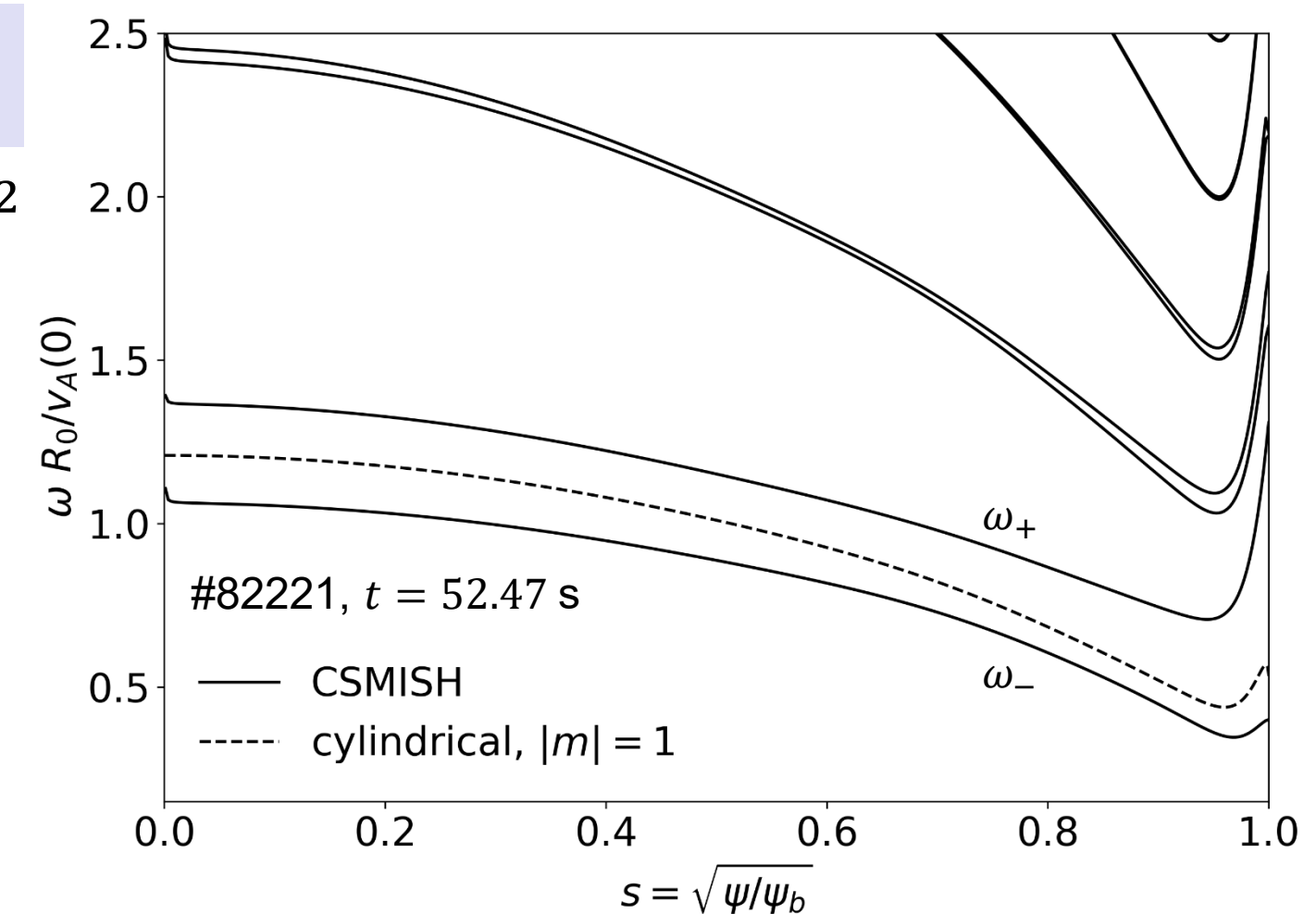
In a first approach, the two lowest branches come from the **coupling of the Fourier harmonics $m = \pm 1 \Rightarrow \xi^A = \xi_{-1}^A e^{-i\theta} + \xi_1^A e^{i\theta}$** . Up to first order terms in κ_s and ε ,

$$[e^{-i\theta} \quad e^{i\theta}] \begin{bmatrix} \frac{(1 + \tilde{\rho} \tilde{\omega}^2 q^2) \kappa_s}{2q^2} & \frac{\tilde{\rho} \tilde{\omega}^2 - \frac{1}{q^2}}{q^2} \\ \frac{\tilde{\rho} \tilde{\omega}^2 - \frac{1}{q^2}}{q^2} & \frac{(1 + \tilde{\rho} \tilde{\omega}^2 q^2) \kappa_s}{2q^2} \end{bmatrix} \begin{bmatrix} \xi_{-1}^A \\ \xi_1^A \end{bmatrix} = 0$$

Equating the determinant to zero yields

$$\tilde{\omega}_{\pm}^2 = \frac{2 \pm \kappa_s}{\tilde{\rho} q^2 (2 \mp \kappa_s)}$$

This analytic formula recovers previous numerical results [12]



$$\varepsilon = \frac{a}{R_0} \begin{cases} a \equiv \text{minor radius} \\ R_0 \equiv \text{magnetic axis position from torus axis} \end{cases}$$

$$\Theta_0(\theta) = 1 + \kappa_s \cos 2\theta + \kappa_a \sin 2\theta$$

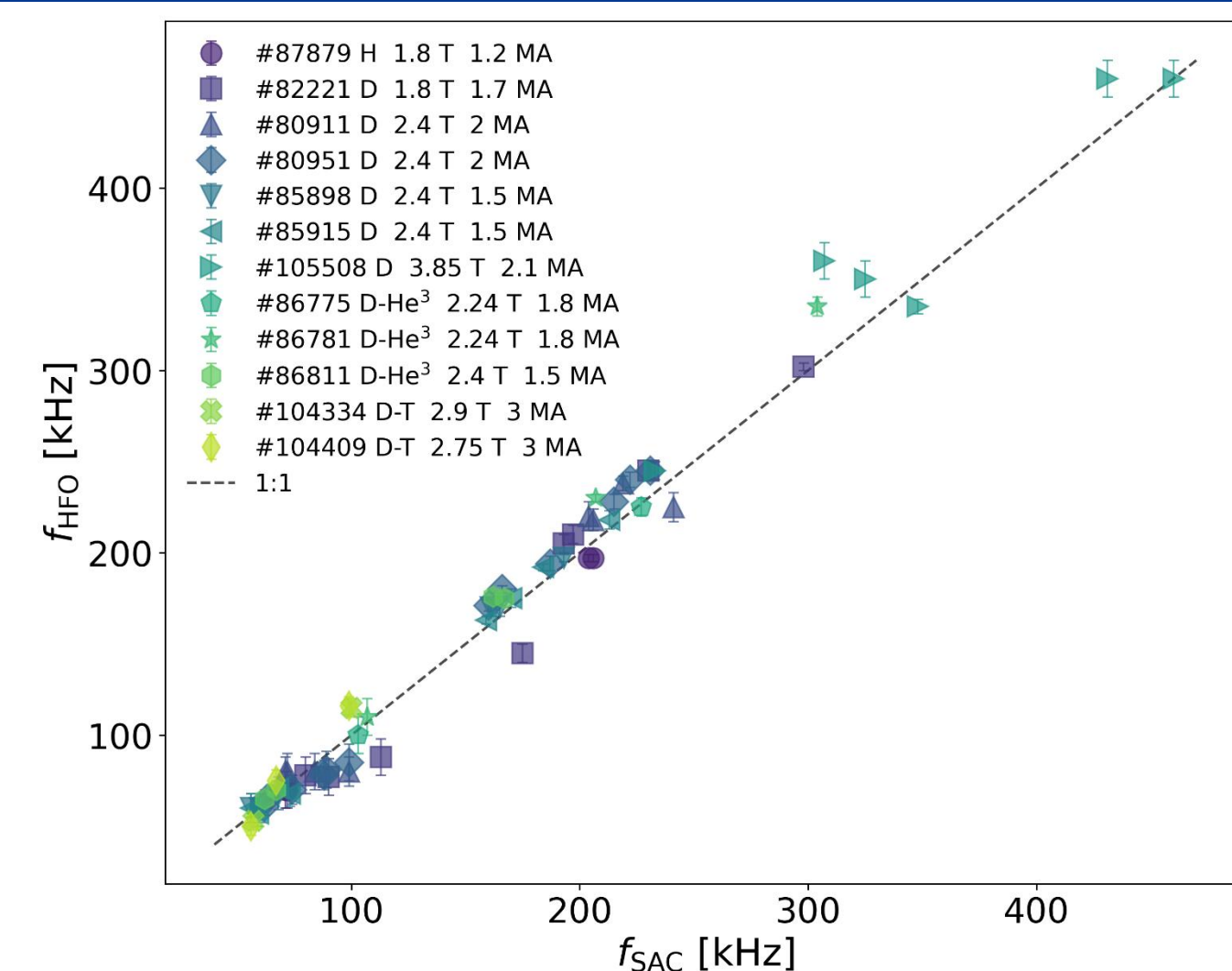
$$\Theta_1(\theta) = \Delta_s \cos \theta + \frac{1}{4} \kappa_a \sin \theta + \eta_s \cos 3\theta + \eta_a \sin 3\theta$$

$$\kappa \approx \sqrt{\frac{1 + \kappa_s}{1 - \kappa_s}} \begin{cases} \kappa \equiv \text{elongation} \\ \kappa_s \equiv \text{ellipticity} \end{cases}$$

$$\xi^A = \xi \cdot \frac{\mathbf{B} \times \nabla \Psi}{|\nabla \Psi|^2} = \sum_m \xi_m^A(\Psi) e^{im\theta}$$

$$\tilde{\omega} = \omega R_0 / v_A(0), \quad \tilde{\rho} = \rho / \rho(0)$$

Experimental measurements versus numerical predictions



The experimentally measured **HFOs frequency (f_{HFO})** agrees with the **frequency of the SAC minima (f_{SAC})** computed by CSMISH.

The analysis was carried out for a comprehensive set of discharge covering a wide set of plasma compositions (H, D, D-He³, D-T), plasma parameters ($I_p = 1.2 - 3.0$ MA, $B = 1.8 - 3.85$ T, $n_e(0) = 2.4 - 6.8$ m⁻³), and heating schemes.

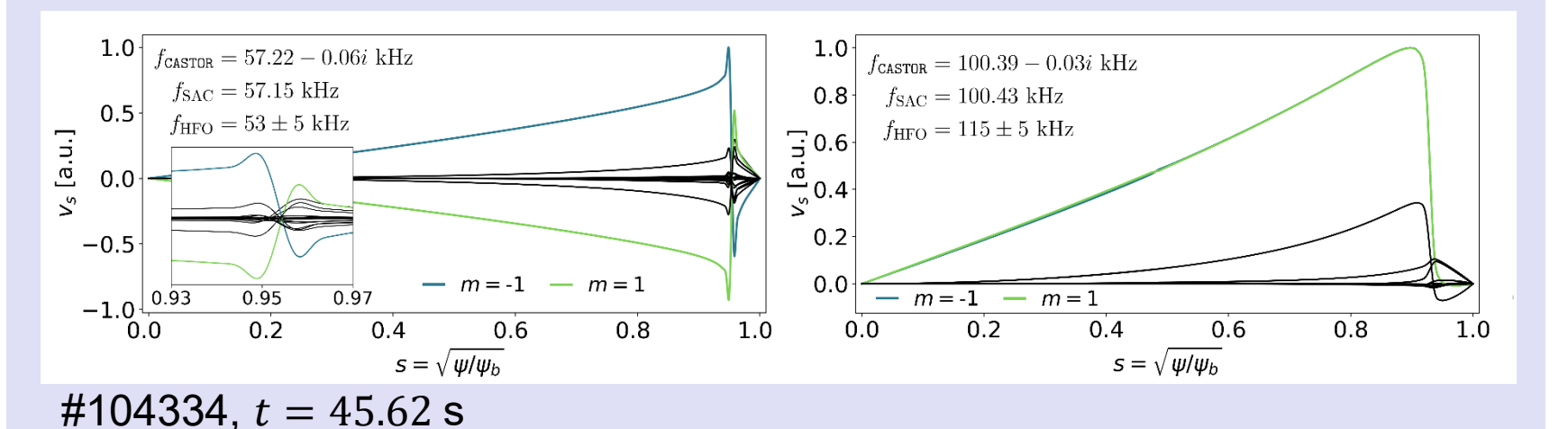
Ideal MHD can accurately predict the frequencies of axisymmetric HFOs

Estimation of GAEs' frequency from the SAC minima is a more **useful** and **robust** approach than mode calculation: MISHKA did not always converge to discrete modes.

Non-ideal terms for numerical GAE structure calculation

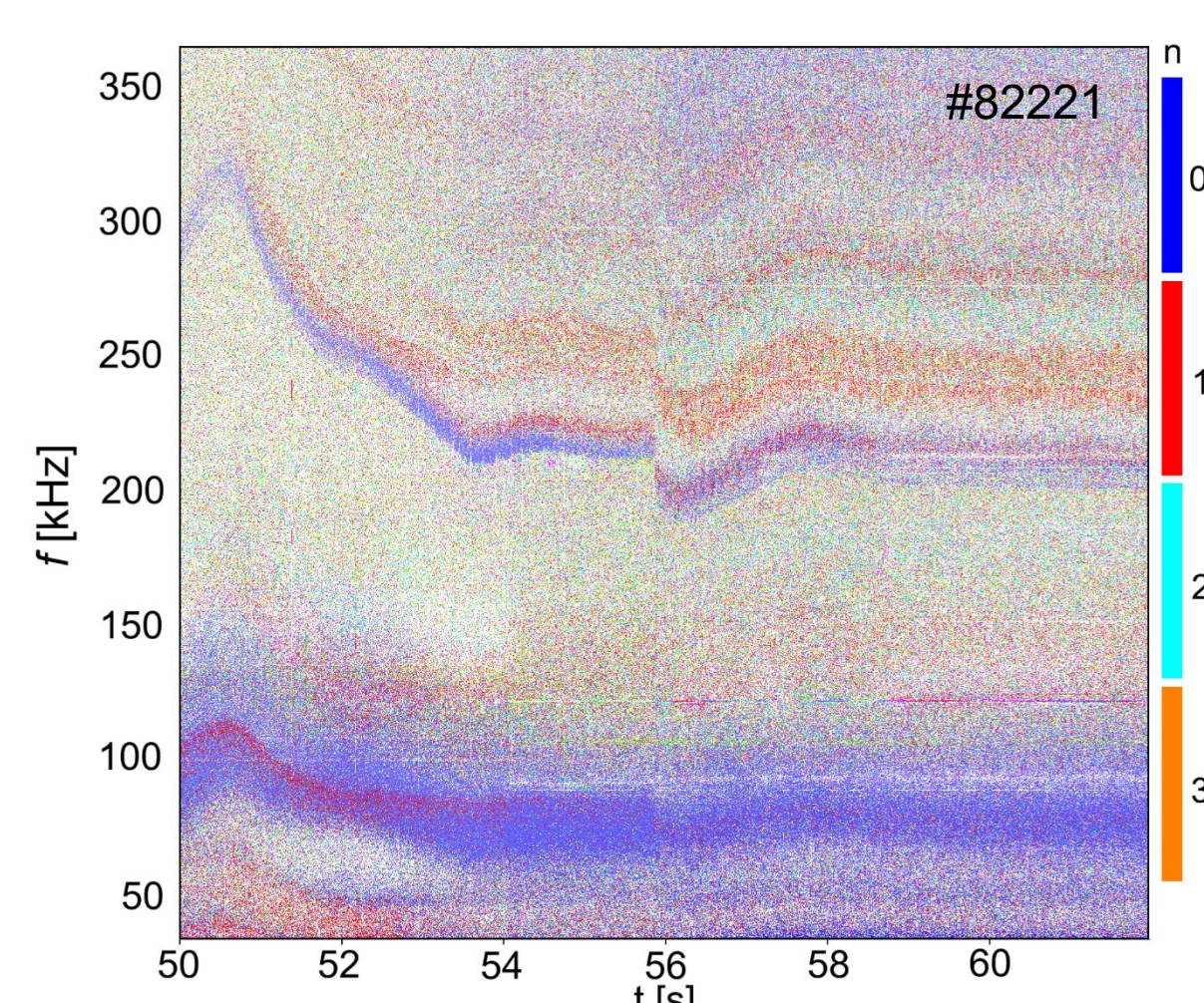
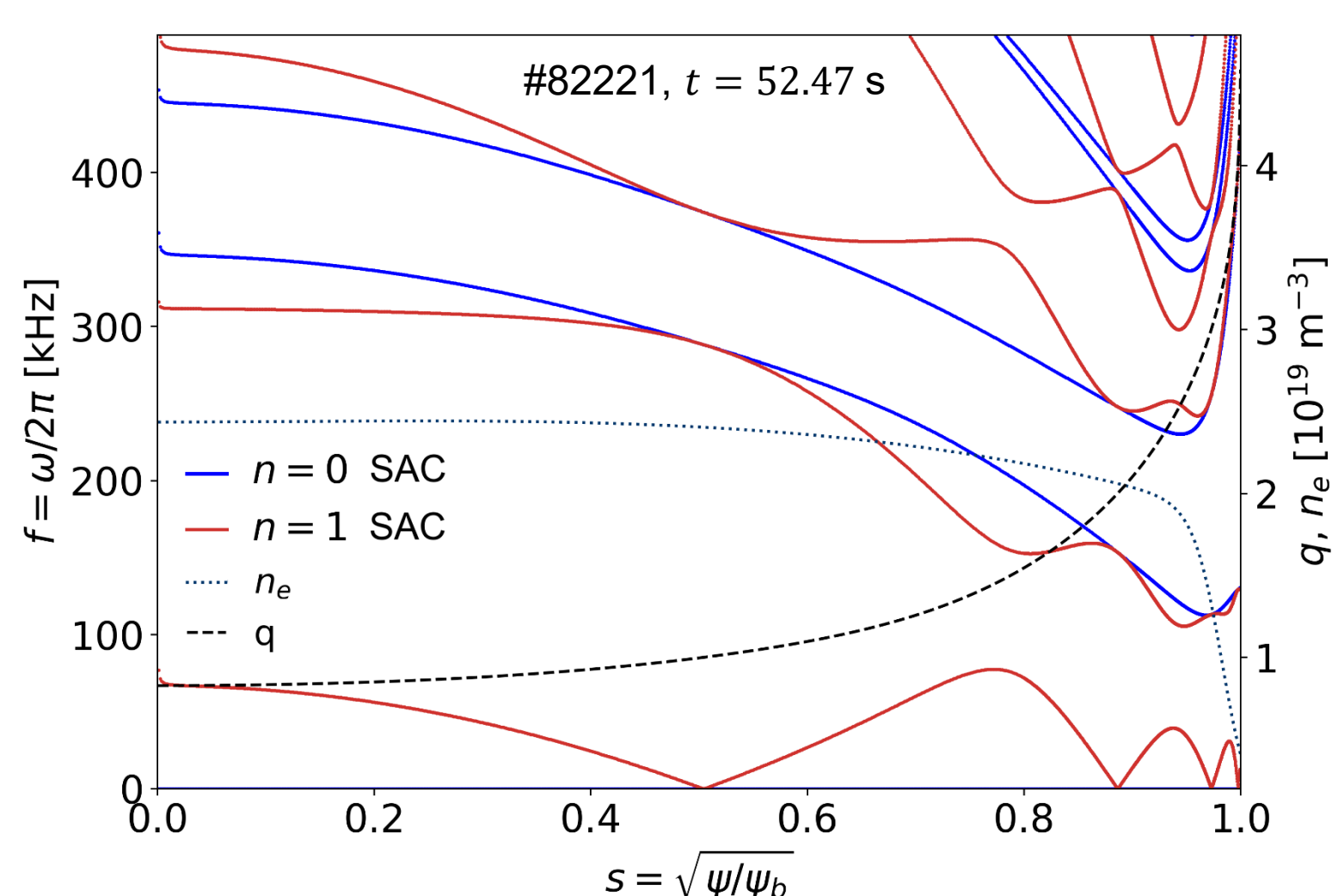
Numerically, as the pedestal grows steeper, the frequency of the GAEs increases, eventually stepping into the SAC $\Rightarrow$ MISHKA converges to singular solutions.

Resistive CASTOR [13] computes global modes with $\Re e(f)$ close to the SAC (from ideal MHD).



Non-axisymmetric HFOs

The frequency modulation and value of $n \neq 0$ HFOs suggest they are **AEs** as well. 1st HFOs: $f_{n=1} \lesssim f_{n=0}$; 2nd HFOs: $f_{n=1} > f_{n=0} \rightarrow$ similar to the SAC branches.



Conclusions and future prospects

Pragmatic interest of HFOs:

- HFOs appear in a variety of plasma scenarios and span a long time window
- f_{HFOs} agree with f_{SAC} at the minima
- Analytical expression for $n = 0$ SAC
 - relates f_{HFOs} with κ_s , q , and ρ

MHD marker for reliable equilibrium reconstruction at the plasma edge

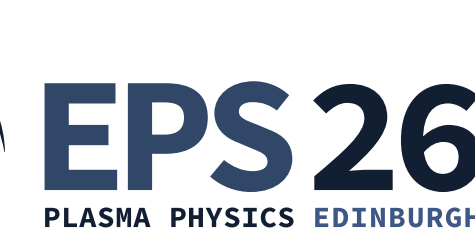
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Acknowledgments

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