

Numerical study of the magnetic field influence on sub-nanosecond laser-foam interaction



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Introduction

This work investigates the influence of **self-generated** and **externally applied** magnetic fields on foam homogenization. The goal is to provide fundamental insights for upcoming experiments to be conducted at the Prague Asterix Laser System (PALS) facility [1].

A new type of low-density foam target is used in our experiments. The foam consists of tightly packed SiO₂ hollow spheres with an average density of *15-40 mg/cc*.

The main motivation for using this foam is its ability to generate long-length-scale plasma for efficient measurement of **parametric instabilities** (SBS, SRS, TPD) under plasma conditions relevant to ICF. In future experiments, the influence of an external magnetic field on parametric instabilities will be studied.

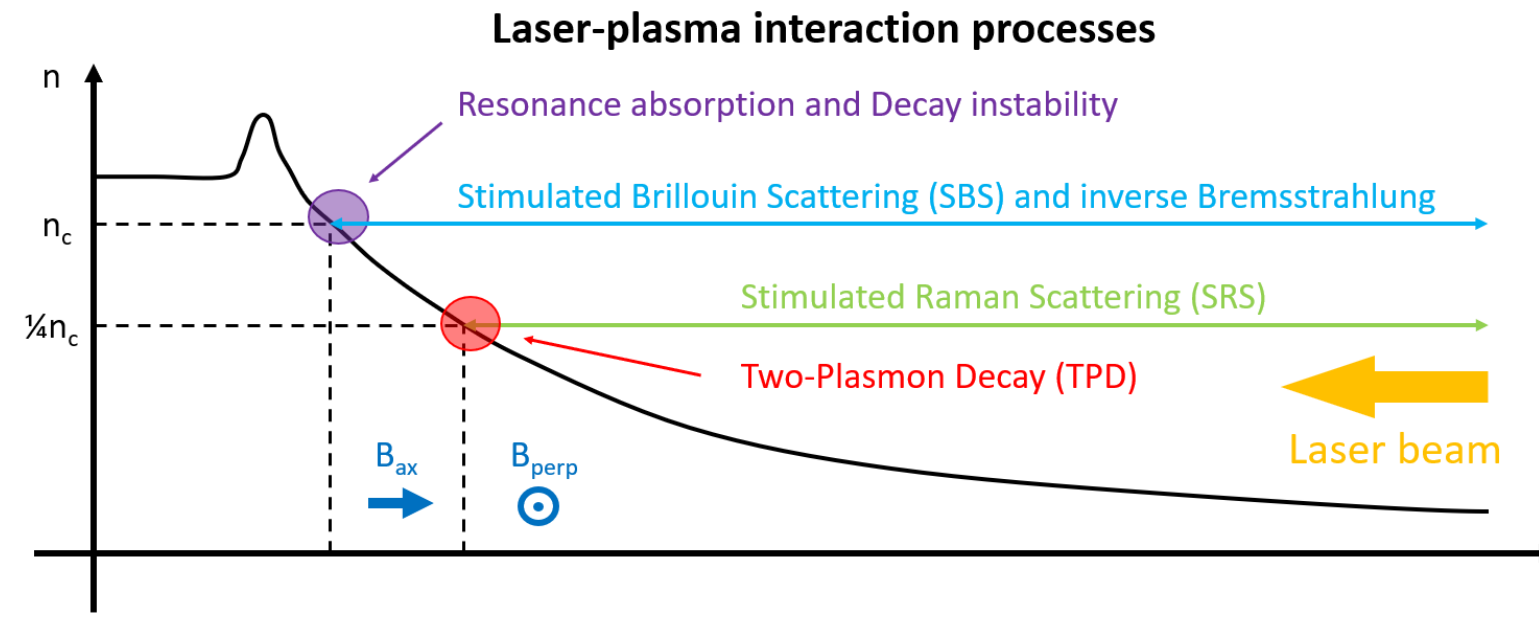
We investigate magnetic field transport during the laser-foam interaction. Due to high laser intensity, the plasma conditions reach a **non-local regime**. Strong electron temperature gradients in the plasma corona lead to magnetic field generation via the **Biermann effect** and to magnetic field advection through the **Nernst effect**. These effects play a dominant role in influencing the macroscopic magnetic field strength and transport. In the non-local regime, however, these terms can be overestimated in the extended MHD model [2].

An **improved limiter** for the heat-flux and thermoelectric terms is proposed, and different models for **anomalous resistivity** are examined.

Simulations are performed using the **FLASH code** [3].

Simulation and experimental setup

- Foam targets with spheres of different diameters – **variable average density**
- Laser beam setup: $E = 200 \text{ J}$, 3ω (438 nm), $50 \mu\text{m}$ focal radius (Gaussian)
- Models used:
 - Radiation hydrodynamics
 - Extended magnetohydrodynamics (extended MHD):
 - Nernst effect** (magnetic field advection)
 - Anisotropic** thermal conductivity κ and magnetic resistivity η
 - Biermann battery** and **Seebeck effect** (self-generated magnetic fields)
 - Hall effect, Cross-field Nernst and Righi-Leduc not included in the simulations
 - FEOS EoS, TOPS opacity tables
 - Cylindrical or Cartesian geometry, laser energy deposition (inverse Bremsstrahlung)



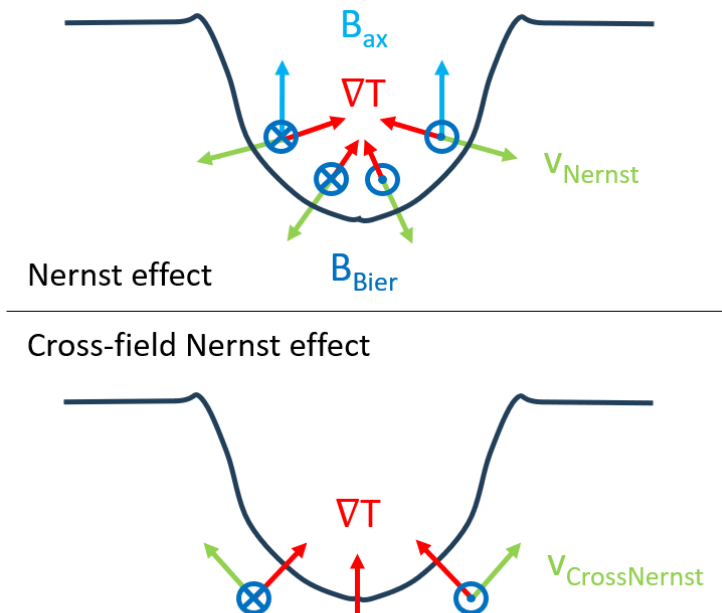
Iodine photodissociation laser at PALS

- Fundamental wavelength: *1315 nm*
- Laser pulse FWHM: *200-350 ps*
- Energy on target: *up to 700 J*
- Laser intensity: *$\sim 3 \times 10^{16} \text{ W/cm}^2$*
- Focus diameter: *$\sim 50-200 \mu\text{m}$*

- External magnetic field applied in parallel (axial) or perpendicular direction to laser beam propagation
- Study the influence of an external magnetic field on parametric instabilities and laser-plasma dynamics
 - Enhanced SBS** backscattering [4]
 - Possible **non-axial part of SRS and TPD inhibition** due to an axial magnetic field [5]
 - Thermal conduction inhibition
- Magnetic field transport under strong temperature gradients
 - Field cavitation** and **compression** due to thermoelectric effects [6]

Magnetic field transport in non-local laser plasma

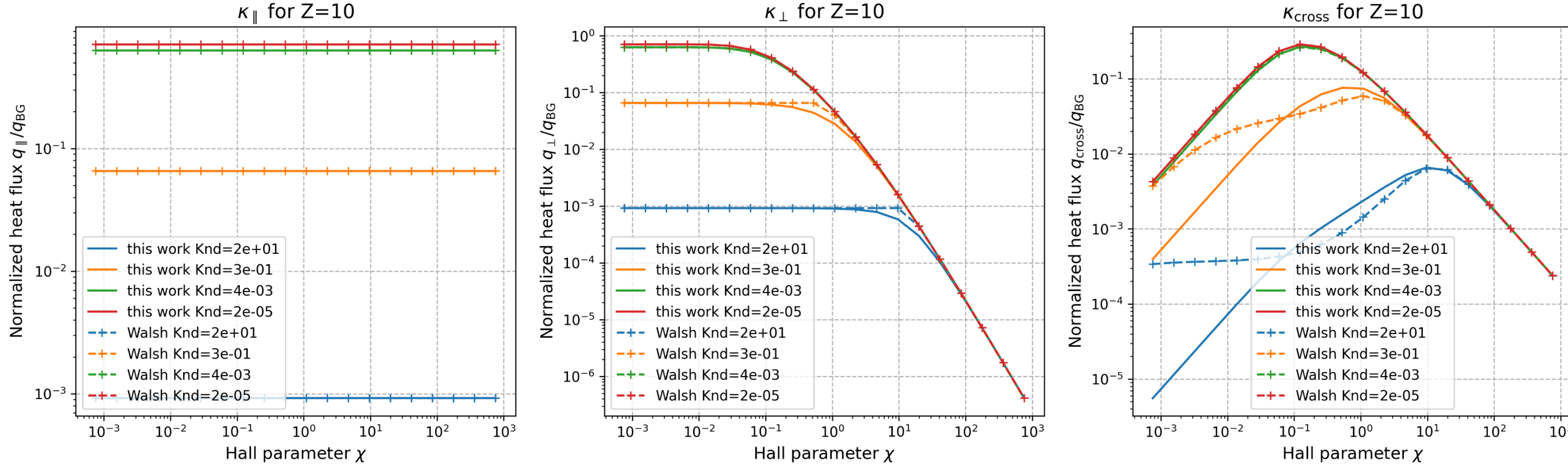
- Strong Nernst advection transports the magnetic field from the hot plasma to the ablation zone behind the critical surface
- Higher magnetization in low-density foam leads to radial magnetic field advection from the shock front (Cross-field Nernst effect)
- External axial magnetic field cavitation and compression behind the critical surface, not advected by Cross-field Nernst at the shock front
- Heat flux needs to be limited in the MHD model due to large temperature gradients



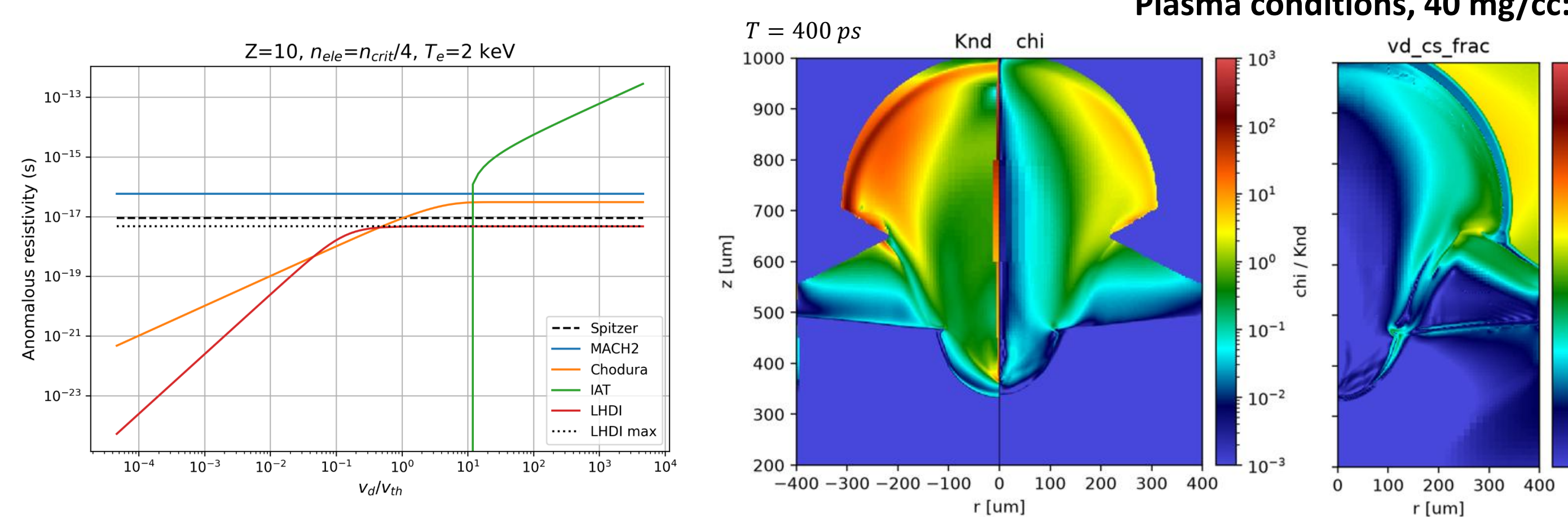
This work:

$$q_{\parallel} = -\frac{\kappa_{\parallel} \nabla_{\parallel} T_e}{1 + \frac{\lambda_d}{f_{new} L_{T_{\parallel}}}}$$
$$q_{\perp} = -\frac{\kappa_{\perp} \nabla_{\perp} T_e}{1 + \frac{\lambda_d}{f_{new} L_{T_{\perp}}} \frac{\kappa_{\perp}^C}{\kappa_{\parallel}^C}}$$
$$q_{\Lambda} = -\frac{\kappa_{\Lambda} \nabla_{\Lambda} T_e}{1 + \frac{\lambda_d}{f_{new} L_{T_{\Lambda}}} \frac{\kappa_{\Lambda}^C}{\kappa_{\parallel}^C}}$$

- Problems:
 - Fixed constant α , no dependence on average ionization
 - Strong magnetization χ (Hall parameter) increases locality – not captured
- Solution:
 - A new form of heat flux limiter based on VFP simulations [7]
 - Non-locality parameter: $\lambda_d = \sqrt{(Z^2 + 0.24Z)/(Z + 4.2)} \lambda_{ei}$, $K_{nd} = \lambda_d / L_T$
 - Limiter for Biermann battery – non-local suppression of $\mathbf{B}$ generation [8]
 - Limiter for thermal electric terms – paired version [7] keeps fractions $\kappa_{\parallel}/\gamma_{\parallel}$, $\kappa_{\perp}/\gamma_{\perp}$ and $\kappa_{\Lambda}/\gamma_{\Lambda}$ **constant**
 - This work uses modified limiters – better scaling of κ_{Λ} for small magnetization [9]



- Classical extended MHD model uses collisional resistivity $\eta_{Spitzer} \sim T_e^{-3/2}$ (with anisotropic coefficients)
- The considered plasma is heated and rapidly expands; large angular magnetic field is generated via Biermann
- Strong currents $\mathbf{j} = \nabla \times \mathbf{B}$ break down the fluid approximation when the electron drift velocity surpasses the kinetic threshold $\rightarrow v_d = |\mathbf{j}|/(en_e) \gtrsim v_i$ or c_s , this leads to anomalous resistivity (scattering on small-scale $\mathbf{E}$ wave fluctuations) $\rightarrow$ micro-instabilities:
- Ion-Acoustic Turbulence** (IAT) – low magnetization, $T_e \gg T_i$, low density
- Lower-Hybrid Drift Instability** (LHDI) – strong magnetization, low density



Conclusions

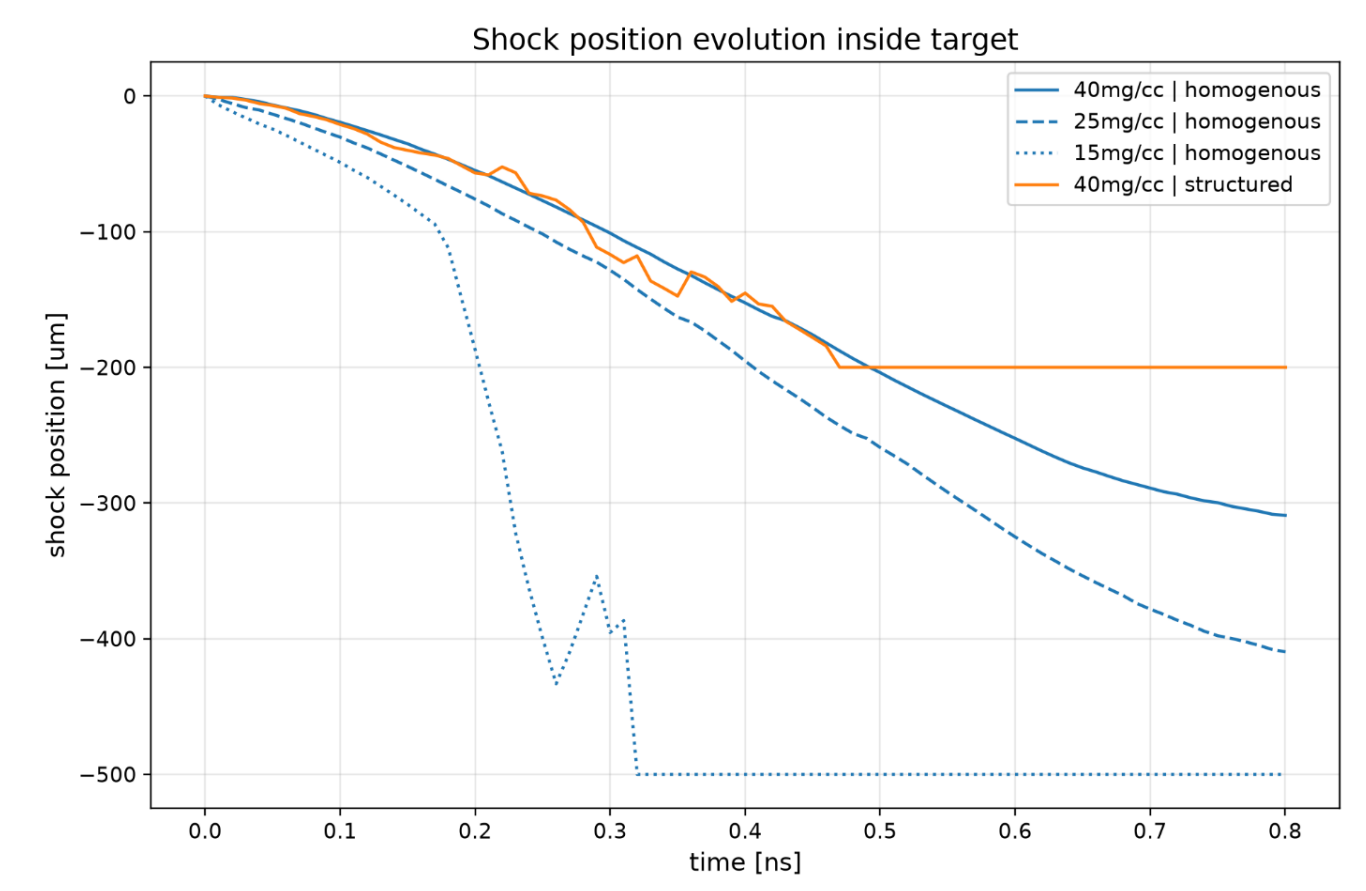
- Simulations of non-local laser plasma conditions using the extended MHD model are difficult – **magnetothermal instability** (Nernst + Righi-Leduc), limiting of ∇T_e terms (Biermann, heat flux, thermoelectric terms), anomalous resistivity – hard to capture kinetic effects accurately in the macroscopic model
- Results show complex magnetic transport can influence shock velocity at subcritical foam densities, leading to fast homogenization
- Improved scaling of anomalous resistivity and Ohm heating fraction (both ion and electron heating through resistivity) is necessary to capture kinetic effects
- A more accurate model of laser beam interaction with foam structure is needed to better resolve shock velocity and homogenization
- Structured simulation of subcritical foam density is needed to evaluate the validity of rapid homogenization

Acknowledgements

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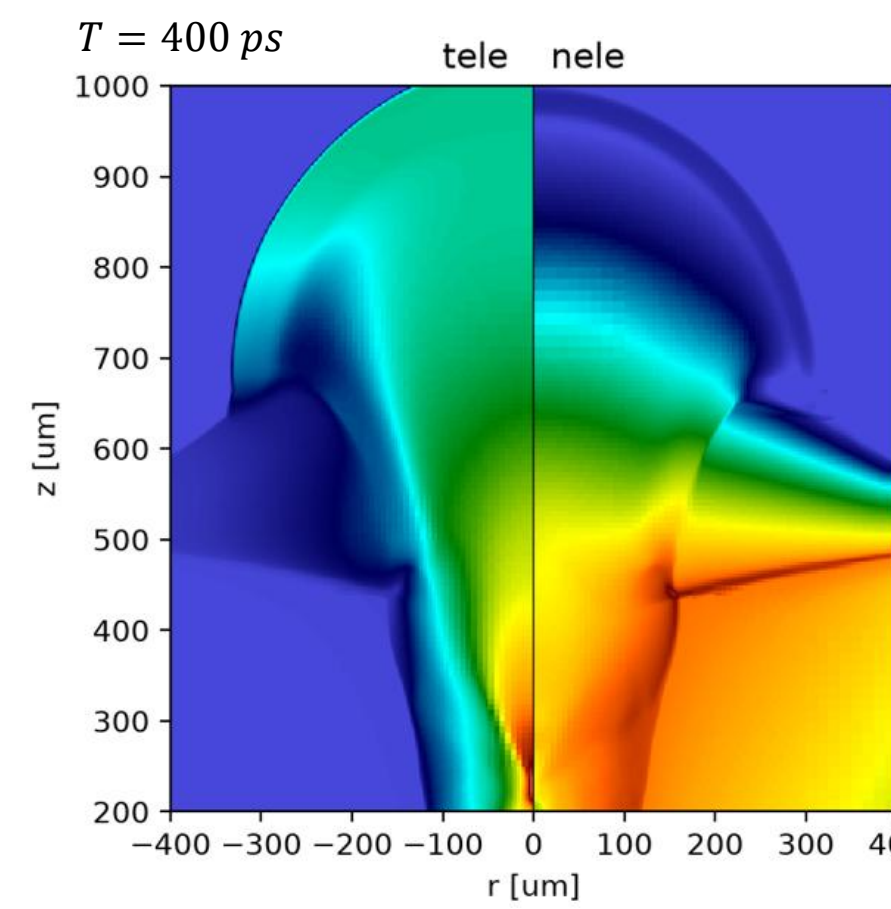
Foam homogenization in external magnetic field

- Shock velocity and laser penetration into the foam target are measured by x-ray streak camera in experiments
- We locate the shock position (depth inside the target) in the extended MHD simulation in the central area of the beam
- There is a clear **dependence of shock velocity on foam density**
- For the subcritical density *15 mg/cc* ($\rho_c \approx 19.3 \text{ mg/cc}$ for 3ω), the simulation shows a **significant jump in shock velocity** at around *180 ps*
- Laser beam penetrates the foam rapidly – 2D density and temperature profiles

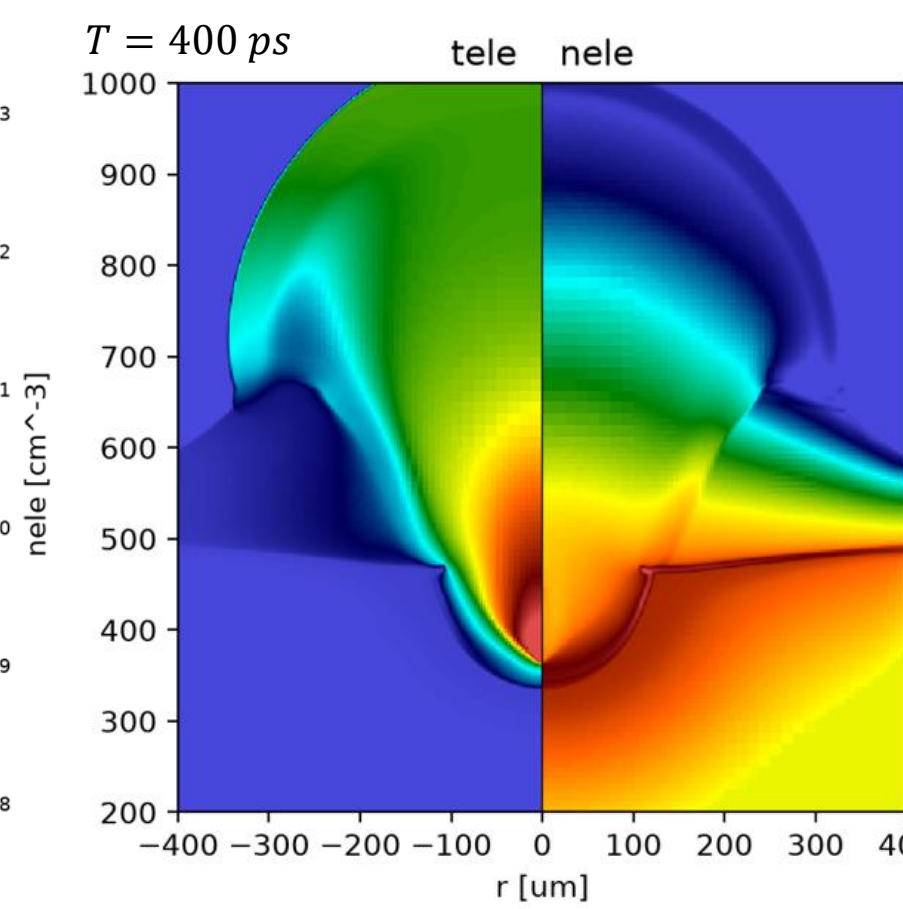


- Shock velocity result from a homogeneous target in the simulation follows the experimental measurement for 3ω
- Highly resolved simulation of foam structure shows agreement with the homogeneous target simulation
- Experimental measurement for 1ω deviates from the simulation $\rightarrow$ possibility of laser beam scattering on the hollow spheres and increased absorption (volume heating) that is not captured by the ray tracing algorithm in the FLASH code

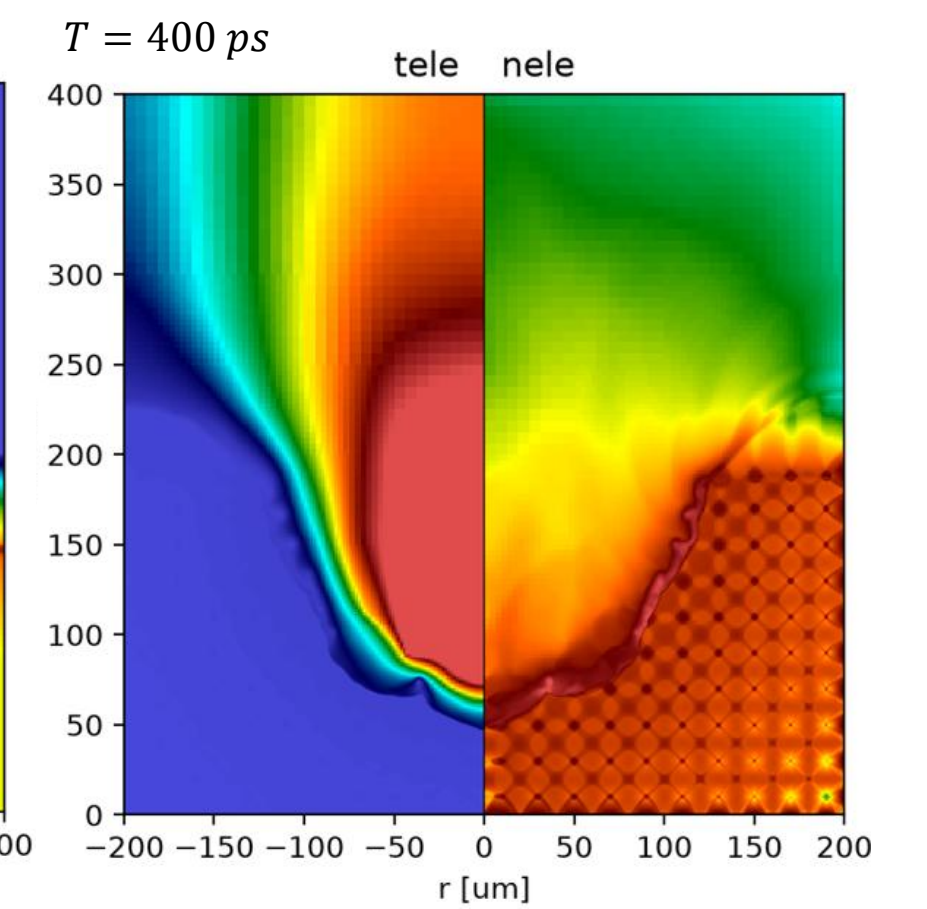
15 mg/cc, homogen.



40 mg/cc, homogen.

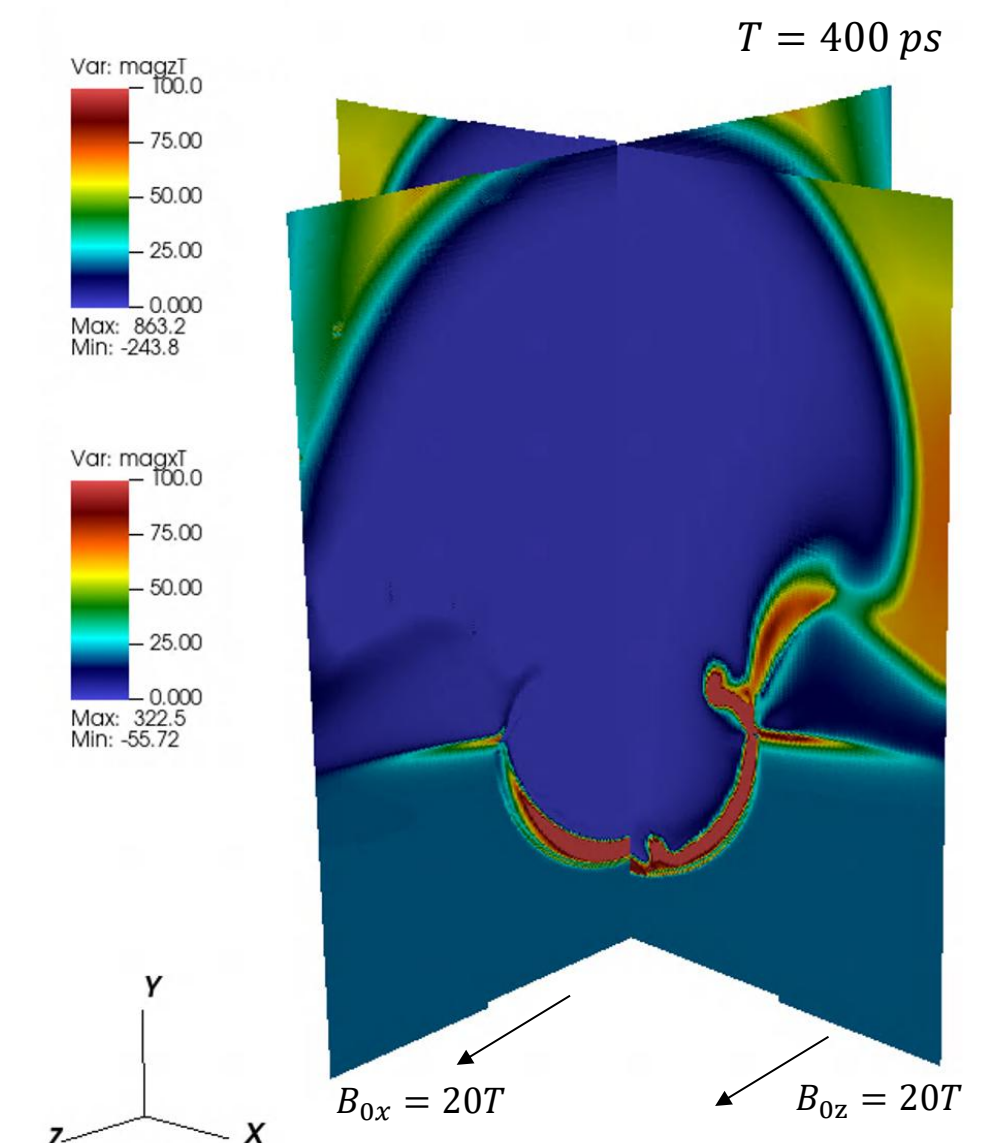
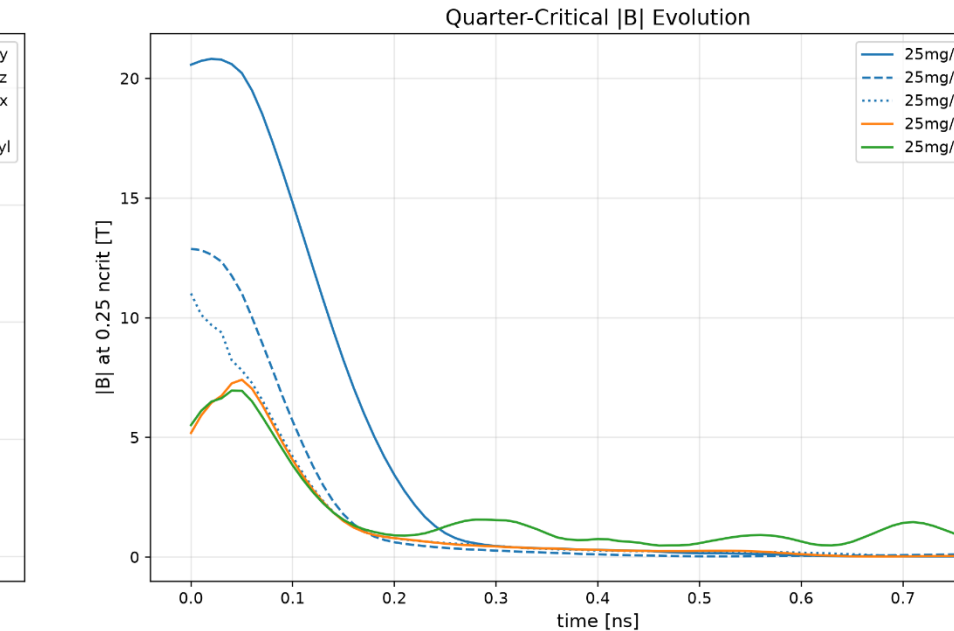
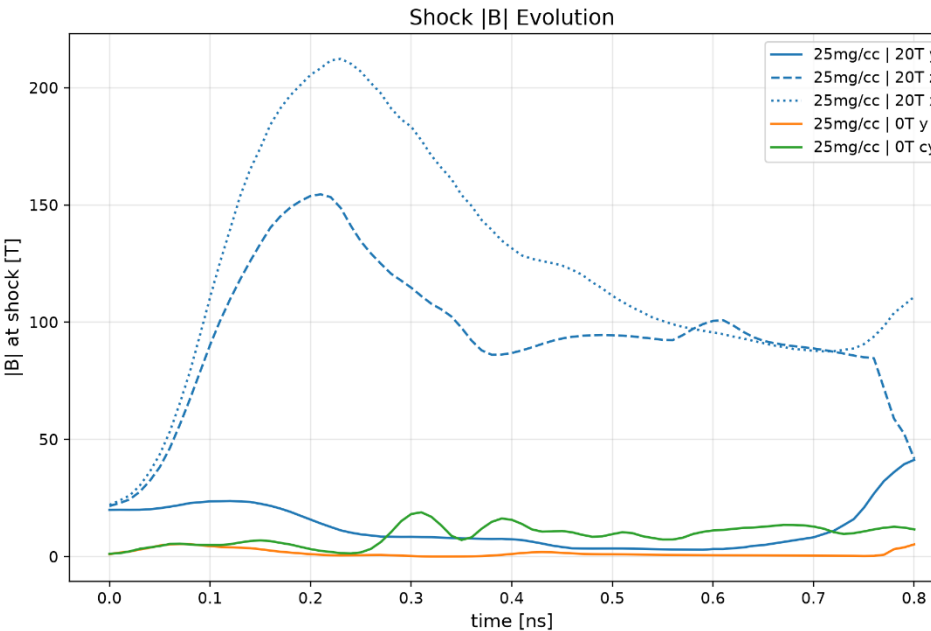


40 mg/cc, structured



- We perform 2D cartesian simulations to examine the influence of different orientations of the initial external magnetic field on homogenization and plasma dynamics
- Two cases: $\mathbf{B} \parallel \text{laser}$, $\mathbf{B} \perp \text{laser}$ (in 2D, this can be either parallel or perpendicular to the simulation domain)
- In all simulations, external and self-generated magnetic fields are strongly advanced from the interaction region via the Nernst effect; at *200 ps*, effectively no magnetic field persists in the heated plasma (for all cases $B_0 = 0, 1, 10, 20 \text{ T}$)
 - The exception is the 2D cartesian simulation with parallel field remaining till *250 ps*

- The magnetic field is advected to the shock front and in the radial direction
 - In both cases of $\mathbf{B} \perp \text{laser}$, the **very strong magnetic field (100-200 T)** is **accumulated** at the shock front compared to the $\mathbf{B} \parallel \text{laser}$ case, where both Biermann and external fields are advected away radially
 - Magnetic field parallel to the computation domain is pushed from the plasma corona, while the perpendicular field is accumulated at the beam edge
 - Although there is no significant difference in shock velocities



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